





Nucleation of cracks and influence of defects in variational phase-field models of brittle fracture

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Fracture mechanics: complex examples



Thin film cracks

Modelling issues:

- Crack nucleation
- Morphogenesis of complex crack patterns

Giant's Causeway

Complex crack patterns using variational phase-field models (Gradient damage models)



Thermal shock, PRL 14



Thin film fracture and debonding, JMPS 14

Full scale simulation in 2d

Ceramic slab from Shao et al. 2011 with $\Delta T = 380 \,^{\circ}\text{C}$

- Material : E = 340 GPa, v = 0.22, $G_c = 42.47 \text{ Jm}^{-2}$, $\sigma_c = 342.2 \text{ MPa}$, $\beta = 8 \times 10^{-6} \text{ K}^{-1}$
- **Dimensions:** $5L \times L \times 0.1L$, with L = 10 mm



A step back: verification and validation on classical test cases

Are these models quantitively accurate?

- Propagation conditions for existing cracks?
- Initiation and nucleation of new cracks?
- Scale effects?
- In linear elastic fracture mechanics, several criteria have been developed to account for experimental findings (e.g. double criterion). Should these criteria be added to variational models ?
- Are energetic criteria enough to predict crack propagation and nucleation?

Outline

- Review of gradient damage models
- Basic examples: propagation of a crack and nucleation in a traction test
- Nucleation at a V-notch
- Nucleation at a U-notch
- Effect of an elliptical hole
- Structural size effects

Sharp vs Smeared approaches to brittle fracture

Sharp models (Griffith)



Cracks as surfaces of sharp discontinuity of the displacement field

Sharp model: Griffith



 $\mathcal{E}(u,\mathcal{K}) = \int_{\Omega\setminus\mathcal{K}} \frac{1}{2} A_0 \varepsilon(u) \cdot \varepsilon(u) \, dx + G_c \operatorname{area}(\mathcal{K})$

Smeared models: gradient damage (phase-field)

• Local minimisation of (α is a irreversible scalar damage field):

$$\{\mathcal{E}_{\ell}(u,\alpha) = \underbrace{\frac{1}{2} \int_{\Omega} (1-\alpha)^2 A \varepsilon(u) \cdot \varepsilon(u) \, dx}_{\text{Strain energy}} + \underbrace{\frac{3G_c}{8} \int_{\Omega} \left(\frac{\alpha}{\ell} + \ell \, \nabla \alpha \cdot \nabla \alpha\right) \, dx}_{\text{Dissipated energy}}\}$$

- Mathematics (Variational regularisations, Γ-convergence): Mumford and Shah, 1989, Ambrosio and Tortorelli, 1990, Bourdin et al. 2000.
- Physics: Phase-field models
 Ginzburg-Landau 1938-1950, Aranson et al., 2000; Karma et al., 2001
- Mechanics: Damage models Local model, non-local models, gradient models, etc ...

Smeared models (Damage)



Cracks approximated by the localisation of a smooth auxiliary field (damage).

Traction of a bar: Griffith model

Assume a crack set \mathcal{K} in the form of *n* transverse cracks and a 1D bar model.



Results for Griffith model

- Local minimality criterion: the bar never cracks
- Global minimality criterion: a single crack at

$$\frac{U_G}{L} = \sqrt{\frac{2G_c}{E_0 L}} \qquad \qquad \sigma_G = \sqrt{\frac{2E_0 G_c}{L}}$$

 Global minimization is the only way to get crack nucleation without adding ad-hoc geometric features (initial test-cracks). But the maximum allowable stresses are infinite for short bar and vanishing for long bars, which is physically not acceptable (see FM98).

Gradient damage model

Total energy functional (linearized isotropic elasticity with isotropic gradient damage)

$$\mathcal{E}_{\ell}(u,\alpha) = \underbrace{\frac{1}{2} \int_{\Omega} A(\alpha) \, \varepsilon(u) \cdot \varepsilon(u) \, dx}_{\text{Strain energy}} + \underbrace{\frac{G_c}{c_w} \int_{\Omega} \left(\frac{w(\alpha)}{\ell} + \ell \, \nabla \alpha \cdot \nabla \alpha \right) dx}_{\text{Dissipated energy}}, \quad c_w = 4 \int_0^1 \sqrt{w(\alpha)} d\alpha$$

- $\mathbf{\alpha}$, a scalar field on Ω : an internal variable representing the damage field.
- $A(\alpha) = a(\alpha) A_0$: the damaged elastic tensor (decreasing from $A(0) = A_0$ to A(1) = 0).
- $w(\alpha)$: dissipation for homogeneous damage (increasing from w(0) = 0).
- l: internal length.

Key requirements

Model retained for computations $a(\alpha) = (1 - \alpha)^2$, $w(\alpha) = \alpha^2 \alpha$

 $c_{w} = 8/3$

- stress softening: $\frac{w'(\alpha)}{s'(\alpha)}$ not increasing, $s = a^{-1}$
- finite elastic limit: w'(0) > 0
- finite dissipation at full damage: w(1) < ∞</p>

Halphen and Nguyen 1975, Fremond, Nedjar 1996; Lorentz, Andrieux 1999, Comi 2001, Nguyen and Andrieux 2005, Benallal, Marigo 2007; Pham, Marigo 2010; Lorentz, Godard 2011; Pham, Marigo 2013

Gradient damage model

Linearized isotropic elasticity with isotropic gradient damage



• Quasi-static time-discrete evolution at time t_i from the state (u_{i-1}, α_{i-1}) :

 $\min_{u \in \mathcal{U}_i, \ \alpha \ge \alpha_{i-1}} \mathcal{E}_{\ell}(u, \alpha), \qquad (\underline{\mathsf{local}} \text{ minimization with unilateral constraint})$

Energetic equivalence with Griffith by setting:

$$c_w = 4 \int_0^1 \sqrt{w(\alpha)} d\alpha \stackrel{\uparrow}{=} \frac{8}{3}$$

Halphen and Nguyen 1975, Fremond, Nedjar 1996; Lorentz, Andrieux 1999, Comi 2001, Nguyen and Andrieux 2005, Benallal, Marigo 2007; Pham, Marigo 2010; Lorentz, Godard 2011; Pham, Marigo 2013

Minimization with bound constraints: variational inequalities

Example: Scalar function of a scalar variable $f : x \in \mathcal{R} \to f(x) \in \mathcal{R}$

$$\frac{\mathbf{x}^*}{x \ge \bar{\mathbf{x}}} = \operatorname*{argmin}_{x \ge \bar{\mathbf{x}}} f(x)$$

First order optimality conditions:

Minimality implies non-negative variations:

$$f(x) - f(\mathbf{x}^*) \simeq \boxed{f'(\mathbf{x}^*)(x - \mathbf{x}^*) \ge 0, \qquad \forall x \ge \bar{\mathbf{x}}}$$

Equivalent KKT conditions:

$$f'(\mathbf{x}^*) \ge 0$$
 $(\mathbf{x}^* - \bar{\mathbf{x}}) \ge 0$ $f'(\mathbf{x}^*)(\mathbf{x}^* - \bar{\mathbf{x}}) = 0$

This is a non-linear problem. If f is convex the solution is unique.





Gradient damage model: evolution problem

First order optimality conditions

Equilibrium equations (u-variation)

div
$$\sigma = 0$$
 in Ω , $\sigma n = f$ on $\partial_f \Omega$, with $\sigma = a(\alpha)A_0\varepsilon(u)$ (1)

Damage criterion and Kuhn-Tucker conditions (α -variation)

$$\ln \Omega \quad : \quad \alpha - \alpha_{i-1} \ge 0, \quad \mathsf{D} \ge 0, \quad (\alpha - \alpha_{i-1}) \,\mathsf{D} = 0, \tag{2}$$

On
$$\partial\Omega$$
 : $\alpha - \alpha_{i-1} \ge 0$, $\frac{\partial\alpha}{\partial n} \ge 0$, $\frac{\partial\alpha}{\partial n} (\alpha - \alpha_{i-1}) = 0$, (3)

where
$$D(\varepsilon, \alpha, \Delta \alpha) = \frac{a'(\alpha)}{2} A_0 \varepsilon \cdot \varepsilon + \frac{G_c}{c_w} \left(\frac{w'(\alpha)}{\ell} - 2\ell \Delta \alpha \right)$$

The energetic criterion (local energy minimality conditions) implies:

Stress domain for homogenous damage:

$$\mathsf{D}(\varepsilon,\alpha,0) \ge 0 \quad \Rightarrow \quad A_0^{-1} \sigma \cdot \sigma \le -\frac{2 \ G_c w'(\alpha)}{c_w \ \ell \ (a^{-1})'(\alpha)}$$

Elastic limit in unixial traction:

$$\sigma_c = \sqrt{\frac{2 G_c E w'(0)}{\ell c_w (a^{-1})'(0)}} \Rightarrow \sigma_c = \sqrt{\frac{3 G_c E}{8\ell}}$$

• Stability (second order condition) depends on the gradient term (and ℓ)

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$$\mathsf{D}(\varepsilon, \alpha, \Delta \alpha) = \frac{\mathsf{a}'(\alpha)}{2} \mathsf{A}_0 \varepsilon \cdot \varepsilon + \frac{\mathsf{G}_c}{\mathsf{c}_w} \left(\frac{w'(\alpha)}{\ell} - 2\ell \Delta \alpha \right)$$

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Considered constitutive models

Homogeneous response in 1d traction



Finite elastic limit

$$\sigma_c = \sigma_M = \sqrt{\frac{3G_cE_0}{8\,\ell}}$$

Linear elastic followed by softening

AT2 - model $w(\alpha) = \alpha^2, \quad a(\alpha) = (1 - \alpha)^2$



- Null elastic limit
- Hardening followed by softening
- Used in most of the papers on phase-field fracture

Pham, Marigo, Maurini 2011; Lorentz, Godard 2011

Propagation, nucleation and internal length

Elementary verification tests

Propagation of an existing crack

Nucleation in a homogenous bar



Illustration: traction of a bar



- Elastic phase (homogeneous damage) ⇒ Loss of stability ⇒ Localized solution (fracture)
- The dissipated energy in the localized solution is equal to the crack length ($G_c = 1$)
- Crack nucleation is a structural stability problem

Propagation of a straight crack

Surfing experience

- Consider a slab with an initial crack I_0
- \blacksquare Impose asymptotic mode-I displacement on the boundary focussed on a crack tip moving at a constant speed v



Top: damage field. Bottom: apply displacement and remove element with lpha > 0.9

Propagation of a straight crack

Evaluate the Energy Release Rate (ERR) through G- θ integral (Destuynder et al., 1983), where θ is an auxiliary field defining a virtual displacement of the domain with



$$G = \int_{\Omega} \left(\frac{a(\alpha)}{2} A_{\mathbf{0}} \varepsilon(u) \cdot \varepsilon(u) \operatorname{div} \theta - a(\alpha) A_{\mathbf{0}} \varepsilon(u) \cdot (\nabla u \nabla \theta) \right) dx$$

Griffith and gradient damage models give the same evolution:

 $G = G_c$, no dependence on ℓ

Propagation of a curved crack



Loading: asymptotic mode-I displacement for propagation along a circular path

Initiation and nucleation tests



Nucleation at a V-notch ?



3- and 4-point bending tests



Experimental results from:

- 4-point bending on Duraluminium by Yosibash et al, IJF 2004
- 3-point bending on PMMA by Dunn at al IJSS 1997 and IJF 1997

Pac-man geometry for the study of the crack nucleation at a notch



Asymptotic solution in linear elasticity

$$\sigma_{\theta\theta} = kr^{\lambda-1}F(\theta), \quad \sigma_{rr} = kr^{\lambda-1}\frac{F''(\theta) + (\lambda+1)F(\theta)}{\lambda(\lambda+1)}, \quad \sigma_{r\theta} = -kr^{\lambda-1}\frac{F'(\theta)}{(\lambda+1)},$$

where $\left| k = \left. \frac{\sigma_{\theta\theta}}{(2\pi r)^{\lambda-1}} \right|_{\theta=0}
ight|_{\theta=0}$ is a generalized stress intensity factor

$$F(\theta) = (2\pi)^{\lambda - 1} \frac{\cos((1+\lambda)\theta) - f(\lambda, \bar{\omega}) \cos((1-\lambda)\theta)}{1 - f(\lambda, \bar{\omega})}, \quad f(\lambda, \bar{\omega}) = \frac{(1+\lambda)\sin((1+\lambda)(\pi-\bar{\omega}))}{(1-\lambda)\sin((1-\lambda)(\pi-\bar{\omega}))}.$$

D. Leguillon and E. Sanchez-Palencia, 1987



Computational method

- Assume separation of scale between inner damage problem and outer linear elastic problem
- Apply asymptotic linear elastic solution at the boundary
- Generalized SIF k is the loading parameter (time)
- $\blacksquare \ \Omega$ large enough, mesh small enough

- The stress are singular for $\bar{\omega} < \pi/2$
- For $\bar{w} = 0$ (a crack), ERR > 0: a crack propagates according to Griffith model.
- For any $\bar{w} > 0$, ERR=0: the stresses are not singular enough to release enough elastic energy for a "Griffith" crack.
- For $\bar{\omega} \simeq \pi/2$ the stress field is almost uniaxial. The critical stress value at nucleation can be estimated by the one of a bar in traction:

$$\sigma_c = \sqrt{\frac{3EG_c}{8\ell}}$$

Damaged crack lips: $\alpha = 1$ on Γ











Critical loading at nucleation





Nucleation at a notch: critical loading at nucleation

Recovering the limit cases : $\bar{w} \rightarrow 0$ (crack) and $\bar{w} \rightarrow \pi/2$ (bar)



AT1 model, Red : damaged crack lips, blue : undamaged crack lips

Comparison with experimental results: V-notch



Critical generalised stress intensity factor at nucleation



U-notch: DENT test



- \blacksquare Consider the regime $a \gg \ell$, $a \gg \rho$
- \blacksquare Study the influence of ρ/ℓ
- Simplify the problem by using asymptotic solution in linear elasticity as a BC

U-notch

Gradient damage vs experimental results for AT1



U-notch



A first summary

Some conclusions from U- and V- notches

- The local energy minimisation principle on the gradient-damage model conciliates the stress and toughness criteria.
- The results are in quantitative agreement with experiments, on many materials and geometries.
- No need for ad-hoc nucleation criteria in this framework.

Influence of imperfections: traction of bar with an hole



 $\sigma_{
m max}\simeq 3\sigma_{\infty}$ for $D\ll W$

 Critical loading of a bar without the hole:

$$U_0(\ell) = L \sqrt{\frac{3G_c}{E8\ell}}, \qquad \sigma_0(\ell) = \sqrt{\frac{3G_cE}{8\ell}}$$

• Two key parameters:
$$\frac{D}{W}$$
, $\frac{\ell}{D}$





(a) Problem setting and loading conditions. (b) Typical mesh used for numerical experiment

■ An analytical solution is available, with maximum stress at A given by

$$\sigma_{\theta\theta} = \sigma\left(1 + \frac{2}{\rho}\right)$$

We apply the asymptotic solution for linear elasticity on the boundary



Critical load for damage nucleation



Critical load for the nucleation of a full crack with $\alpha = 1$ on a band



- I (small hole) $\sigma = \sigma_0$: The effect of the hole is filtered out in the relative large processing zone. Nucleation as in a homogeneous bar.
- **II** $(a \sim \ell)$ Non-trivial transition zone with the material scale effect
- III (large hole) $\sigma = \sigma_0/3$: Nucleation at end of the elastic phase.

Critical load for the crack nucleation: shape effect



• For $\rho = 0.1$ and $a \gg \ell$ we recover the scaling law of a crack:

$$\sigma \propto \sqrt{\frac{EG_c}{a}}$$

Crack in a domain of finite width and scale effects



■ $a \ll \ell$ Stress criterion (no effect of the defect)

- $a \gg \ell$ Griffith criterion (no effect of the the internal length)
- $a \simeq \ell$ "Mixed" regime

Conclusions

- Quantitively accurate results for crack nucleation on different geometries.
- V-notch and U-notch experimental results can be used to accurately estimate the parameters of the models.
- No need for additional local ad-hoc nucleation criterion, the energy minimisation principle implies both.
- This is a nonlinear damage model, not only a criterion for the limit of validity of linear elasticity: can predict post-critical behaviour.
- Crack nucleation is a structural stability problem.





Movies





- FEniCS/python based, parallel open-source demo codes in the form of interactive on-linejupyter-notebook available here (no installation required) https://notebooks.azure.com/cmaurini/libraries/varfrac
- Scripts reproducing most of the results of this presentations are available (thanks to Tianyi Li and Erwan Tanné): https://bitbucket.org/cmaurini/gradient-damage
- See also Blaise Bourdin's fortran90 code (https://bitbucket.org/bourdin/mef90-sieve)
- Tanné, E., Li, T., Bourdin, B., Marigo, J.-J., Maurini, C. (2018). Crack nucleation in variational phase-field models of brittle fracture. Journal of the Mechanics and Physics of Solids, 110, 80–99. https://doi.org/10.1016/j.jmps.2017.09.006